# Electromagnetic wave propagation in complex materials with thermo-electric coupling

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First Interdisciplinary Workshop,

of the German-Russian Interdisciplinary Science Center (G-RISC)

**Structure and Dynamics of Matter** 

October 18 — 20, 2010, Berlin, Germany

#### 1. Coupled system of Maxwell's equations and heat equation

Let  $\Omega \subset \mathbb{R}^3, \partial \Omega$  the smooth enough boundary,



Fig. 1

$$\varepsilon E_t + \sigma E = \operatorname{rot} H,$$
  

$$\mu H_t + \operatorname{rot} E = 0,$$
  

$$\theta_t - \Delta \theta = \sigma |E|^2,$$
(1)

where E(x,t) is the electric field, H(x,t) is the magnetic field,  $\theta(x,t)$  is the temperature,  $\sigma(\theta)$  is the electrical conductivity,  $\varepsilon(x)$ is the electric permittivity,  $\mu(x)$  is the magnetic permeability,  $(x,t) \in Q_T := \Omega \times [0,T)$ .

$$E(x,t) = (E_1(x,t), E_2(x,t), E_3(x,t), x = (x_1, x_2, x_3),$$

$$\mathsf{rotE} = \begin{vmatrix} i & j & k \\ \frac{\partial}{\partial x_1} & \frac{\partial}{\partial x_2} & \frac{\partial}{\partial x_3} \\ E_1 & E_2 & E_3 \end{vmatrix} = \nabla \times \mathsf{E}$$

 $q = \sigma |E|^2$  is Joule's heat



Fig. 2

Initial and boundary conditions (Landau/Lifshetz, Electrodynamics):

 $\nu \times E(x,t) = \nu \times G(x,t)$ , where  $\nu$  is the normal to the boundary  $\partial \Omega$ , G(x,t) is some function.

$$eta_{
u} = 
u \cdot 
abla heta = 0, \, (x,t) \in S_T := \partial \Omega \times [0,T],$$
  
 $E(x,0) = E_0(x), H(x,0) = H_0(x), heta(x,0) = heta_0(x), x \in \Omega,$ 

 $E_0(x), H_0(x)$ , and  $\theta_0(x)$  are some functions.

#### 2. The one-dimensional case

Suppose that  $Q_T = \{(x,t) | 0 < x < 1, 0 < t < T\}$  and E(x,t) = (0, e(x,t), 0), H(x,t) = (0, 0, h(x,t)).



Fig. 3

Rewrite the previous system in the form

$$\begin{aligned} \varepsilon e_t + \sigma e &= -e_2 h_x, \\ \mu h_t + e_x &= 0, \\ \theta_t - \theta_{xx} &= \sigma |e|^2, \quad x \in (0, 1), t > 0, \\ e(0, t) &= f(t), \theta(0, t) = 0, \quad t > 0, \\ e(x, 0) &= e_0(x), h(x, 0) = h_0(x), \quad x \in (0, 1), \\ \theta(x, 0) &= \theta_0(x). \end{aligned}$$

$$(2)$$

We assume that  $w(x,t) := \int_{0}^{t} e(x,s)ds$  and  $\varepsilon(x) = \mu(x) = 1$ and derive the following system:

$$\begin{cases} w_{tt} = w_{xx} - \sigma(\theta)w_t, \\ \theta_t - \theta_{xx} = \sigma(\theta)w_t^2, \\ w(0,t) = f_1(t), w(1,t) = f_2(t), \\ \theta(0,t) = \theta(1,t) = 0, \\ w(x,0) = w_0(x), w_t(x,0) = w_1(x), \\ \theta(x,0) = \theta_0(x). \end{cases}$$
(3)



Fig. 4

#### Assumptions

- (A1) 1)  $\sigma(z)$  is piecewise  $C^1$  on  $(0, \infty)$ .
  - 2) There exist constants  $0 < \sigma_0 \leq \sigma_1$ , such that  $\sigma_0 \leq \sigma(z) \leq \sigma_1$ , for any  $z \geq 0$ .
  - 3) There exists a constant  $\sigma_2$ , such that  $\sigma(z)z \leq \sigma_2$ , for any  $z \geq 0$ .
- (A2) 1)  $\theta_0(x)$  is nonnegative and of class  $C^2(0, 1)$ .
  - 2)  $w_0(x), w_1(x)$  are of class  $C^4(0, 1)$  and  $f_1(t), f_2(t)$  are of class  $C^2$ .

## Remark 1

Hölder spaces (Ladyzhenskaya, Solonnikov, Ural'ceva, 1968) Suppose  $\Omega$  is a bounded region in  $\mathbb{R}^n$ .

$$C^{\alpha}\left(\bar{\Omega}\right) = \left\{ f \in C_{b}\left(\bar{\Omega}\right) : \sup_{t,s\in\bar{\Omega},t\neq s} \frac{|f\left(t\right) - f\left(s\right)|}{|t-s|^{\alpha}} < +\infty \right\}$$

$$C^{\alpha,0}\left(\bar{\Omega}\times[a,b]\right) = \left\{ f \in C\left(\bar{\Omega}\times[a,b]\right) : f\left(\cdot,t\right)\in C^{\alpha}\left(\bar{\Omega}\right)\forall t\in[a,b] \right\}$$

$$C^{0,\alpha}\left(\bar{\Omega}\times[a,b]\right) = \left\{ f \in C\left(\bar{\Omega}\times[a,b]\right) : f\left(x,\cdot\right)\in C^{\alpha}\left[a,b\right]\forall x\in\bar{\Omega} \right\}$$

$$C^{\alpha,\frac{\alpha}{2}}\left(\bar{\Omega}\times[a,b]\right) = C^{0,\frac{\alpha}{2}}\left(\bar{\Omega}\times[a,b]\right) \cap C^{\alpha,0}\left(\bar{\Omega}\times[a,b]\right)$$

$$C^{2+\alpha,1+\frac{\alpha}{2}}\left(\bar{\Omega}\times[a,b]\right) = \left\{ f \in C^{2,1}\left(\bar{\Omega}\times[a,b]\right) : f_{t},f_{x_{i}x_{j}}\in C^{\alpha,\frac{\alpha}{2}}\left(\bar{\Omega}\times[a,b]\right) \right\}.$$

Under the assumptions (A1)-(A2), the system has a classical unique global solution(Yin)  $(w, w_t, \theta)$  on  $Q_T = (0, 1) \times (0, T)$ , for any  $T < \infty$ . Furthermore,  $w(x, t) \in C^{3,3}(\bar{Q}_T), \theta(x, t) \in C^{2+\alpha, 1+\frac{\alpha}{2}}(\bar{Q}_T)$  for some  $\alpha \in (0, 1)$ .

#### 3. Almost periodic perturbations

Let  $f_1(t)$  and  $f_2(t)$  be are almost periodic functions.

**Definition 1** A scalar continuous function  $f : \mathbb{R} \to \mathbb{R}$ , is called *almost-periodic in the sense of Bohr*, if for every  $\varepsilon > 0$  there exists an  $L = L(\varepsilon)$  such that every interval  $[t_0, t_0 + L]$  contains at least one number  $\alpha$  for which  $|f(t + \alpha) - f(t)| < \varepsilon, t \in \mathbb{R}$ .

#### Example 1

- a) periodic functions:  $\cos t$ ,  $\sin t$  etc.
- b)  $a_1 \cos \omega_1 t + a_2 \cos \omega_2 t, \omega_1/\omega_2$  irrational,  $a_1 \neq 0, a_2 \neq 0$

#### 4. Cocycles in infinite-dimensional phase space

**Definition 2** Let  $(P, \rho)$  be a complete metric space,  $\{\tau^t\}_{t \in \mathbb{R}} : P \to P$  is a continuous map satisfying

1) 
$$\tau^{0}(\cdot) = id_{P}$$
,

2)  $\tau^{t+s}(\cdot) = \tau^t \circ \tau^s, \forall t, s \in \mathbb{R}.$ 

Then  $(\{\tau^t\}_{t\in\mathbb{R}}, P)$  is called *base flow.* 

#### **Example 2**

Shift map  $\tau^t(p) := p + t, \ \forall t, p \in \mathbb{R} = P$ 

**Definition 3** Let *X* be a complete metric space. A *cocycle over the base flow*  $\{\tau^t\}_{t\in\mathbb{R}}$  is a family of maps  $\{\varphi^t(p,\cdot): X \to X\}_{t\in\mathbb{R}, p\in P}$  such that

1) 
$$\varphi^0(p,\cdot) = id_X, \forall p \in P,$$

2) 
$$\varphi^{t+s}(p,\cdot) = \varphi^t(\tau^s(p),\varphi^s(p,\cdot)), \forall t,s \in \mathbb{R}, \forall p \in P.$$



#### 5. Constructing a cocycle for the microwave heating problem

Introduce the auxiliary functions  $f(x,t) = f_1(t)(x-1) + f_2(t)x$ ,  $x \in (0,1), t \ge 0, W(x,t) = w(x,t) - f(x,t), V(x,t) =$  $W_t(x,t) - f_t(x,t)$ . Rewrite equation (3) as the first order system

$$\begin{cases} W_t = V - f_t, \\ V_t = W_{xx} - \sigma(\theta)V + f_{tt}, \\ \theta_t - \theta_{xx} = \sigma(\theta)V^2, \\ W(0,t) = W(1,t) = 0, \theta(0,t) = \theta(1,t) = 0, \\ W(x,0) = w_0(x) - f(x,0), \\ W_t(x,0) = w_1(x) - f_t(x,0), \\ \theta(x,0) = \theta_0(x). \end{cases}$$
(4)

Then (4) can be formally written as system

$$\frac{du}{dt} = Au + Bg(V,\theta) + F(t), \tag{5}$$

where  $u = (W, V, \theta)$ ,  $F(t) = (-f_t, f_{tt}, 0)$ , A, B are linear operators.

If  $(W(x,t), V(x,t), \theta(x,t))$  is a solution of (4), we can write it as

$$u(t, t_0, u_0) = (W(\cdot, t), V(\cdot, t), \theta(\cdot, t)),$$
  

$$u_0 = (w_0(x) - f(x, 0), w_1(x) - f_t(x, 0), \theta_0(x)),$$

Consider the shift map  $\tau^t : \mathbb{R}_+ \to \mathbb{R}_+$  and the family of mappings  $\{\varphi^t(t_0, \cdot)\}_{t, t_0 \in \mathbb{R}_+}$  such that

$$\varphi^{(\cdot)}(\cdot, \cdot) : \mathbb{R}_+ \times \mathbb{R}_+ \times X \to X,$$
  

$$\tau^t(t_0) = t + t_0,$$
  

$$\varphi^t(t_0, u_0) = u(t + t_0, t_0, u_0),$$

for any  $t \in \mathbb{R}_+, t_0 \in \mathbb{R}_+, u_0 \in X$ , where X is the Banach space  $X = H_0^1(0, 1) \times L^2(0, 1) \times L^2(0, 1)$  with the norm  $\| (W, V, \theta) \|_X^2 = \| W \|_{H_0^1}^2 + \| V \|_{L^2}^2 + \| \theta \|_{L^2}^2$ .

## 6. Existence of an almost periodic solution

**Definition 4** A continuous function  $y : \mathbb{R}_+ \to X$  is called an *integral* of the cocycle on  $\mathbb{R}_+$ , if  $\varphi^t(s, y(s)) = y(t + s)$  for all  $t, s \in \mathbb{R}_+$ .

The integral  $y : \mathbb{R}_+ \to X$  of the cocycle is said to be

1) *uniformly stable* if for any  $\varepsilon > 0$ , there exist  $\delta = \delta(\varepsilon) > 0$  such that

$$\varphi^t(s, N_{\delta}(y(s))) \subset N_{\varepsilon}(y(t+s)), \ t, s \in \mathbb{R}_+,$$

2) *uniformly asymptotically stable* (UAS) if it is uniformly stable and there exist a  $\delta_0 > 0$  with the property that for any  $\varepsilon > 0$ , there is a  $t_0 > 0$  such that

$$\varphi^t(s, N_{\delta_0}(y(s))) \subset N_{\varepsilon}(y(t+s)), \ t \ge t_0, s \in \mathbb{R}_+.$$

**Theorem 1** System (2) has an almost periodic solution which is uniformly asymptotically stable.

Idea of the proof: We consider the functional

$$\Phi(W, V, \theta) = \int_{0}^{1} (W_x^2 + 2\lambda WV + V^2 + a\theta^2) dx,$$

where  $\lambda > 0$  and a > 0 are some parameters and show that each solution of (4) is bounded in future, uniformly asymptotically stable and has a compact orbit.

Consider the system (4) with initial-boundary conditions

$$W(x,0) = p \sin(\pi x),$$
  

$$V(x,0) = 0, \theta(x,0) = 0,$$
  

$$W(0,t) = \sin(t) + \sin(\sqrt{2}t), W(1,t) = 0,$$
  

$$\theta(0,t) = \theta(1,t) = 0.$$



Fig. 6

#### Reference

[1] Yu. Kalinin, V. Reitmann and N. Yumaguzin, Asymptotic behavior of Maxwell's equation in one-space dimension with thermal effect, submitted for publication in AIMS Journals.

## 7. Existence of the global attractor for the 1-dimensional microwave heating problem

We consider the system

$$w_{tt} - w_{xx} + \sigma(\theta) w_t = 0, \quad 0 < x < 1, \quad t > 0 \theta_t - \theta_{xx} = \sigma(\theta) w_t^2, \qquad 0 < x < 1, \quad t > 0$$
(6)

with boundary conditions (defining the set  $\equiv$ )

$$w(0,t) = w(1,t) = 0, \quad t > 0 \theta(0,t) = \theta(1,t) = 0, \quad t > 0$$
(7)

and initial conditions

$$w(x,0) = w_0(x), w_t(x,0) = w_1(x), \theta(x,0) = \theta_0(x), \quad (8)$$
  
  $x \in (0,1)$ 

We make the following assumptions:

**(A3)** 
$$w_0 \in C^3[0,1], w_1 \in C^3[0,1], \theta_0 \in C^{2+\alpha}[0,1]$$

(A4) Compatibility conditions (defining the set 
$$\Lambda$$
)  
 $\theta_0(0) = \theta_0(1) = 0, w_0(0) = w_0(1) = 0,$   
 $w_1(0) = w_1(1) = 0$   
 $\theta_0''(0) = \theta_0''(1) = 0, w_0''(0) = w_0''(1) = 0$ 

(A5)  $\sigma$  is piecewise  $C^1$  on  $[0, +\infty)$  and there exists M > 0 and  $k \ge 0$  such that  $0 \le \sigma(z) \le M(1+z)^k, \forall z \ge 0$ 

**Theorem 2** (Yin, 2001) Under the assumptions (A3) – (A5) the system (6) - (8) has a unique global solution  $(w, \theta)$  with  $\theta \in C^{2+\alpha,1+\frac{\alpha}{2}}(\bar{Q_T}), w \in C^{3,3}(\bar{Q_T})$  for any T > 0, where  $Q_T = (0,1) \times (0,T]$ .

#### Idea of the proof: Define the phase space

$$X = (C^{3}[0,1] \times C^{2}[0,1] \times C^{2+\alpha}[0,1]) \cap \Lambda \cap \Xi$$

The norm is induced from  $H_0^1 \times L^2 \times H_0^1$ 

 $\|(w, v, \theta)\|_X^2 = \|w_x\|_{L^2}^2 + \|v\|_{L^2}^2 + \|\theta_x\|_{L^2}^2, \text{ where}$  $\|u\|_{L^2} = (\int_0^1 u^2 dx)^{1/2}.$ 

Define the map  $\varphi^t : X \to X$ :

$$\varphi^t u_0 = u(t) = u(t, u_0),$$

where  $u(t) = (w(t), v(t), \theta(t))$  is the solution with initial data  $u_0 = (w_0, w_1, \theta_0)$ 

 $\sigma : \mathbb{R}_+ \to \mathbb{R}$  is called *Lipschitz* if  $\exists L > 0 \ \forall x, y \in \mathbb{R}_+$ :  $|\sigma(x) - \sigma(y)| \le L|x - y|$  11

**Theorem 3** Suppose additionally that  $\sigma$  is Lipschitz and  $0 < \sigma_0 \leq \sigma(z) \leq \sigma_1, \forall z \geq 0$ . Then the mapping  $\varphi^t : X \to X$  is continuous for any  $t \geq 0$ .

**Proof**  $u = (w(t), v(t), \theta(t))$ ,  $\overline{u} = (\overline{w}(t), \overline{v}(t), \overline{\theta}(t))$  are solutions with initial data  $u_0 = (w_0, v_0, \theta_0)$ ,  $\overline{u}_0 = (\overline{w}_0, \overline{v}_0, \overline{\theta}_0)$ .

$$\begin{split} \tilde{w} &:= w - \bar{w} \\ \tilde{\theta} &:= \theta - \bar{\theta} \\ \tilde{w}_{tt} - \tilde{w}_{xx} + \left(\sigma\left(\theta\right)v - \sigma\left(\bar{\theta}\right)\bar{v}\right) = 0 \ \left(\cdot, \tilde{w}_{t}\right) \\ \tilde{\theta}_{t} - \tilde{\theta}_{xx} &= \sigma\left(\theta\right)v^{2} - \sigma\left(\bar{\theta}\right)\bar{v}^{2} \ \left(\cdot, \tilde{\theta}_{t}\right) \\ 1)\frac{1}{2}\frac{d}{dt}\left(\|\tilde{w}_{x}\|^{2} + \|\tilde{w}_{t}\|^{2}\right) \leq C\left(\left\|\tilde{\theta}\right\|^{2} + \|\tilde{w}_{t}\|^{2}\right) \\ 2)\frac{1}{2}\frac{d}{dt}\|\tilde{\theta}_{x}\|^{2} + \|\tilde{\theta}_{t}\|^{2} \leq C\|\tilde{\theta}\|^{2} + C\|\tilde{w}_{t}\|^{2} \\ \frac{d}{dt}\Phi\left(t\right) \leq C\Phi\left(t\right), \text{ and } \Phi\left(t\right) = \|\tilde{w}_{x}\|^{2} + \|\tilde{w}_{t}\|^{2} + \|\tilde{\theta}_{x}\|^{2} \\ \Phi\left(t\right) \leq \Phi\left(0\right)e^{ct} \end{split}$$

C denotes different constants that may depend on  $T, \|u_0\|, \|\overline{u}_0\|$ 

**Theorem 4** The initial-boundary problem (6) - (8) defines a dynamical system on X.

#### 8. Absorbing set

**Definition 5** The set  $A \subset X$  is called a *global attractor (B-attractor)* for the dynamical system  $(\{\varphi^t\}, X)$  if it is compact, invariant and attracts all points (bounded sets) of the space X.

**Definition 6** The set  $Z \subset X$  is called *absorbing* for the dynamical system  $(\{\varphi^t\}, X)$  if for any bounded that  $M \subset X$  there exists a  $t_0 = t_0(M)$  such that  $\varphi^t M \subset Z, \forall t \ge 0$ .



Fig. 7

**Idea of the proof:** Find a smooth functional  $\Phi$  on X such that

$$\frac{d}{dt}\Phi\left(\varphi^{t}u\right) + c_{1}\Phi\left(\varphi^{t}u\right) \le D_{1}, \forall t > 0$$
$$\Phi\left(u\right) \ge c_{2}||u||^{2} - D_{2}, \forall u \in X$$
$$u, c_{2}, D_{1}, D_{2} > 0.$$

where  $c_1, c_2, D_1,$ 

#### **Additional assumptions**

(A6) 
$$\sigma(\theta) \theta \leq \sigma_2$$
 for  $\theta < R$   
(A7)  $\exists \sigma_0, \sigma_1 0 < \sigma_0 \leq \sigma(\theta) \leq \sigma_1$   
(A8) There exist  $0 < \lambda < 1$  and  $a > 0$  such that

$$4\lambda\sigma_0 - 4\lambda^2 - 2\lambda a\sigma_2 - \lambda^2\sigma_1^2 > 0$$

 $||u||_{L^{\infty}} = \operatorname{ess\,sup}\{|u(x)| : x \in (0,1)\}$ 

**Theorem 5** The solution components of (6) – (8)  $w_x$  and  $w_t$  decay to zero exponentially in  $L^2(0,1)$ ,  $\theta$  decays to zero in  $L^{\infty}(0,1)$  with rate  $Ct^{-\frac{1}{2}}$  and exponentially in  $L^2(0,1)$ , when  $\|\theta\|_{L^{\infty}} < R$ .

$$||u(t)||_{L^2} \le C_1 e^{-C_2 t} ||u(0)||, \forall t \ge 0$$

**Corollary 1** Each neighborhood of zero is an absorbing set for (6) -(8).

## Idea of the proof:

Consider the functional on X

$$\Phi(w,v,\theta) = \int_0^1 \left( w_x^2 + 2\lambda wv + v^2 + a\theta^2 \right) dx$$

Differentiating  $\Phi$  with respect to the system we obtain (using the Poincare inequality and integration by parts):

$$\frac{d}{dt}\Phi(t) = \int_0^1 \left(-\lambda w_x^2 - \lambda\sigma(\theta) wv + (\lambda - \sigma(\theta) + a\sigma(\theta)\theta)v^2 - a\theta_x^2\right) dx \le -C_3\Phi$$

**Theorem 6** Let the assumptions 1), 2) be satisfied for  $\|\theta\|_{L_2} > M$ . Then the ball { $(w, v, \theta : \Phi(w, v, \theta) < M$ } is an absorbing set for the system (6) – (8). If additionally we have  $\|\theta\|_{L_{\infty}} < R$ , the attraction rate is exponential.

## 9. Use of the frequency domain method

A similar system was considered by Likhtarnikov and Yakubovich (1977):

$$w_{tt} + 2\varepsilon w_t - \Delta w + \alpha w = f(\theta) ,$$
  

$$\theta_t - \beta \Delta \theta + w - \gamma g(\theta) = 0 ,$$
  

$$w(x, 0) = w_0(x) , w_t(x, 0) = w_1(x) , \theta(x, 0) = \theta_0(x) .$$

Nonlinearities f and g satisfy the following conditions:

$$\theta g(\theta) - f^{2}(\theta) \ge 0, \exists G : G(\theta) \ge 0 : G'(\theta) = g(\theta)$$

Consider the equation

$$\dot{y} = Ay + B\xi, w = Cy.$$
 (9)

We introduce the spaces:

 $Y_0$  is a Hilbert space

 $Y_1 = D(A)$  is a Hilbert space dense in  $Y_0$ , A is a linear self-adjoint operator,

 $Y_{-1}$  is anti-dual to  $Y_1$  ( $Y_{-1}$  is the closure of  $Y_0$  with respect to the scalar product  $(\eta, \varrho)_{-1} = (A^{-1}\eta, A^{-1}\varrho)_0$ ),

 $Y_1 \subset Y_0 \subset Y_{-1}$  is a Hilbert space rigging structure,

 $\equiv$  and W are Hilbert spaces,

B, C are linear operators  $\Xi \to Y_{-1}, Y_0 \to W$ ,

 $F^{j}(y,\xi)$  are Hermitian forms on  $Y_1 \times \Xi$ ,  $j = 1, \ldots, m$ ,

$$F^{j}(y,\xi) = (F_{1}^{j}y,y)_{-1,1} + 2\operatorname{Re}(F_{2}^{j}y,\xi)_{\Xi} + (F_{3}^{j}\xi,\xi)_{\Xi},$$

"Control":  $\xi = \varphi(y)$ 

Consider the set  $N = \bigcap_{1}^{m} N_{j}$  of processes  $(y(\cdot), \xi(\cdot))$ , satisfying the constraints

 $F^{j}\left(y\left(t
ight),\xi\left(t
ight)
ight)\geq0$ 

(local constraint) or

$$\int_{0}^{t_{k}}F^{j}\left(y\left(t\right),\xi\left(t\right)\right)dt\geq\gamma_{j}$$

(integral constraint)



**Fig. 8** 

$$(\Pi_0(i\omega)\xi,\xi) = \sum_{j=1}^m \tau_j F^j(\chi(i\omega)\xi,\xi) + \delta C\chi(i\omega) \|\xi\|^2,$$

where  $\chi(i\omega) = (i\omega I - A)^{-1}B$  is the frequency characteristic of the system (9).

**Definition 7** The system (9) is called *absolutely W*-*dichotomic* if for any process  $p \in (L + p_0) \cap N$ , the state of which is bounded  $||x(t)||_0 \leq C_0 \quad \forall t \geq 0$ , we have

$$w(\cdot) \in L^{2}(0, +\infty; W), \|w(\cdot)\|_{2,W}^{2} \leq C\left(\|x_{0}\|_{0}^{2} + \sum \gamma_{i}\right).$$
(10)

**Definition 8** The system (9) is called *absolutely W*-stable, if for any process  $p \in (L + p_0) \cap N$  we have (10).

**Theorem 7** (The quadratic criterium of absolute stability of Likhtarnikov and Yakubovich)

Let the following conditions hold:

1) The pair (A, B) is  $L^2$ -controllable;

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- 2) A has no spectral points in some neighborhood of the imaginary axis ;
- 3) The system (9) is minimally stable in the class N.

Then the system (9) is absolutely stable if and only if the frequency domain-condition is satisfied:

 $\begin{aligned} \exists \delta > 0 \quad \forall \omega \in \mathbb{R} \ \forall \xi \in \Xi : \\ F\left(\chi\left(i\omega\right)\xi,\xi\right) \leq -\delta \|\xi\|^2. \end{aligned}$ 

#### 10. Application of the frequency-domain method

Define the system

$$\dot{y} = Ay + B\xi,$$

$$y(x,t) = \begin{pmatrix} y_1 \\ y_2 \\ y_3 \end{pmatrix} = \begin{pmatrix} w_t(x,t) \\ w(x,t) \\ \theta(x,t) \end{pmatrix},$$

$$A = \begin{pmatrix} -\sigma_0 I & -A_0 & 0 \\ I & 0 & 0 \\ 0 & 0 & -A_0 \end{pmatrix}, B = \begin{pmatrix} I & 0 \\ 0 & 0 \\ 0 & I \end{pmatrix}$$

 $A_0$  is the operator defined by  $-\Delta$  with homogeneous boundary conditions.

#### **Notations:**

$$\sigma(\theta) = \sigma_0 + \hat{\sigma}(\theta), \sigma_0 > 0$$
  
$$\xi(x, t) = \varphi(y) = \begin{pmatrix} -\sigma(\hat{\theta})w_t(x, t) \\ (\sigma_0 + \hat{\sigma}(\theta))w_t^2(x, t) \end{pmatrix}$$

 $\{\lambda_k\}_{k=1}^{\infty}$  are the eigenvalues of  $A_0$ ,  $0 < \lambda_1 \le \lambda_2 \le \dots \lambda_k \to +\infty$ ,  $\{e_k\}_{k=1}^{\infty}$  are the eigenfunctions which define an orthonormized basis in  $L^2$ .

 $w^{k}(t), \theta^{k}(t), \xi^{k}(t)$  are the Fourier coefficients of functions w(x, t),  $\theta(x, t), \xi(x, t)$ , respectively with respect to  $\{e_k\}$ .

$$ilde{u}(i\omega) = \int_0^\infty e^{-i\omega t} u(t) {
m d} t$$

Decompose solution components by the basis and take the Fourier transform in t (k = 1, 2, ...):

$$-\omega^2 \tilde{w}^k(i\omega) + i\omega \sigma_0 \tilde{w}^k(i\omega) - \lambda_k \tilde{w}^k(i\omega) = \tilde{\xi}_1^k(i\omega) ,$$
  
$$i\omega \tilde{\theta}^k(i\omega) + \lambda_k \tilde{\theta}^k(i\omega) = \tilde{\xi}_2^k(i\omega) .$$

From here we obtain:

$$\tilde{w}^{k}(i\omega) = \tilde{\xi}_{1}^{k}(i\omega)\chi_{0}(i\omega,\lambda_{k}), \tilde{\theta}^{k}(i\omega) = \tilde{\xi}_{2}^{k}(i\omega)\chi_{1}(i\omega,\lambda_{k}),$$

where

$$\chi_0(i\omega,\lambda_k) = \left(-\omega^2 + \sigma_0 i\omega + \lambda_k\right)^{-1}, \chi_1(i\omega,\lambda_k) = (i\omega + \lambda_k)^{-1},$$

$$\left(\Pi_{0}\left(i\omega\right)\tilde{\xi},\tilde{\xi}\right)=\sum_{k}\left(\Pi_{0}^{k}\left(i\omega\right)\tilde{\xi}^{k},\tilde{\xi}^{k}\right)\,.$$

Let us consider two quadratic constraints:

1) 
$$F^{1}(y,\xi) = \int_{0}^{1} y_{1}\xi_{1}dx = \int_{0}^{1} w_{t}(x,t)\widehat{\sigma}(\theta)dx \ge 0$$
,  
 $\left(\prod_{0}^{k}(i\omega)\,\tilde{\xi}^{k},\tilde{\xi}^{k}\right) = \operatorname{Re}\left(-\overline{\tilde{\xi}}_{1}^{\mathsf{k}}\widetilde{\mathsf{W}}_{\mathsf{t}}^{\mathsf{k}}\right) = \frac{-\sigma_{0}\omega^{2}}{(\lambda_{\mathsf{k}}-\omega^{2})^{2}+\sigma_{0}^{2}\omega^{2}}|\overline{\tilde{\xi}}_{1}^{\mathsf{k}}|^{2}.$ 

$$\Pi_{0}^{k}(i\omega) = \begin{pmatrix} \frac{-\sigma_{0}\omega^{2}}{(\lambda_{k}-\omega^{2})^{2}+\sigma_{0}^{2}\omega^{2}} & 0\\ 0 & 0 \end{pmatrix}$$

Frequency-domain condition:  $\Pi_0^k(i\omega) \leq 0, \ \forall \omega \in \mathbb{R} \text{ if } \sigma_0 > 0.$ 

Counterexample: blow-up if  $\varphi_2(y) = \theta^3$  $\theta_t - \theta_{xx} = \theta^3$  on (0, 1) with Dirichlet BC.



Fig. 9

2)  $F^2(y,\xi) = \int_0^1 (\theta^2 - a\theta\xi_2) dx = \int_0^1 (\theta^2 - a\theta\sigma(\theta) w_t^2) dx$ Additional assumption :

$$\int_{0}^{1} \left(\theta^{2} - a\theta\sigma\left(\theta\right)w_{t}^{2}\right)dx > 0$$

$$F = F^{1} + F^{2}$$

$$\Pi_{0}^{k}\left(i\omega\right) = \begin{pmatrix} \frac{-\sigma_{0}\omega^{2}}{\left(\lambda_{k} - \omega^{2}\right)^{2} + \sigma_{0}^{2}\omega^{2}} & 0, \\ 0 & \frac{1 - a\lambda_{k}}{\omega^{2} + \lambda_{k}^{2}} \end{pmatrix}$$

$$a := \frac{2}{\lambda_{1}}$$

Frequency-domain condition:  $\Pi_0^k(i\omega) \leq 0, \ \forall \omega \in \mathbb{R}.$ 

Theorem 7: Absolute stability with respect to the outputs  $w, w_t, \theta$ .

## **Experimental results**

1)
$$\sigma(\theta) = \begin{cases} \exp(8-\theta), \theta > 4; \\ \exp(\theta), \theta \le 4 \end{cases}$$
  
 $\theta_0(x) = 9(1 - |2x - 1|)$   
a)  $w_0(x) = 5(1 - |2x - 1|), w_1(x) = 0$ 



b) The same  $\sigma$  and  $\theta_0$ ,

 $w_0(x) = 0, w_1(x) = 5(1 - |2x - 1|)$ 



Fig. 10

2) 
$$\sigma(\theta) = \begin{cases} \theta - 3, \theta > 4; \\ 5 - \theta, \theta \le 4 \end{cases}$$
  
 $\theta_0(x) = 9(1 - |2x - 1|)$   
 $w_0(x) = 5(1 - |2x - 1|)$   
 $w_1(x) = 0$ 



Fig. 11

## 11. Absolute stability in the microwave heating problem

Let us consider the following system:

$$\begin{split} w_t &= v, & x \in (0,1), \ t > 0, \ (11) \\ v_t &= w_{xx} - \sigma(\theta)v, & x \in (0,1), \ t > 0, \ (12) \\ k(\theta)_t &= \theta_{xx} + \sigma(\theta)v^2, & x \in (0,1), \ t > 0, \ (13) \\ w(0,t) &= 0, \ w(1,t) = 0, & t > 0, \ (14) \\ \theta(0,t) &= \theta(1,t) = 0, & t > 0, \ (15) \\ w(x,0) &= w_0(x), \ v(x,0) = w_1(x), & x \in (0,1), \ (16) \\ \theta(x,0) &= \theta_0(x), & x \in (0,1). \ (17) \end{split}$$

In equation (13),  $k(\theta)$  is a function defining internal energy of the material at the given temperature  $\theta$ . With respect to Joule's law and heat capacity definition we have

$$k(\theta)_t = \frac{\partial k}{\partial \theta} \theta_t = c(\theta) \theta_t,$$

where c is the heat capacity of the material.

Thus, equation (13) takes the classic form of the heat transfer equation

$$c\theta_t = \theta_{xx} + F.$$

## Assumptions

- (A9) 1)  $\sigma(z)$  is piecewise  $C^1$  on  $(0, \infty)$ , k(z) is a smooth function on  $(0, \infty)$ .
  - 2) There exist  $\sigma_0$ ,  $\sigma_1$ :  $0 < \sigma_0 \le \sigma_1$ , such that  $\sigma_0 \le \sigma(z) \le \sigma_1$ , for  $z \ge 0$ .
  - 3) There exists  $k_1 > 0$ , such that  $k_1 \le k'(z)$ , for  $z \ge 0$ .
  - 4) Either exists  $\sigma_2$ , such that  $\sigma(z)z \leq \sigma_2$ , or  $k_2$ , such that  $k'(z)^{-1}z \leq k_2$ , for  $z \geq 0$ .
- (A10) 1) θ<sub>0</sub>(x) is nonnegative and of class C<sup>2</sup>(0,1). Compatibility conditions up to the second order hold at the points (0,0) and (1,0).
  2) w<sub>0</sub>(x) and w<sub>1</sub>(x) are of class C<sup>4</sup>(0,1). Compatibility holds at (0,0) and (1,0).

(Yin): Under conditions (A9) – (A10), the system (11)-(17) has a unique classical solution  $u = (w, v, \theta)$  on  $Q_T = (0, 1) \times$ (0, T), for  $T < \infty$ . Furthermore,  $w(x, t) \in C^{3,3}(\bar{Q}_T), \ \theta(x, t) \in C^{2+\alpha, 1+\frac{\alpha}{2}}(\bar{Q}_T)$  for some  $\alpha \in (0, 1)$ .

Consider the family of mappings  $\{\varphi^t(\cdot)\}_{t\in\mathbb{R}_+}, \varphi^{(\cdot)}(\cdot) : \mathbb{R}_+ \times X \to X$  of the form

$$\varphi^t(u_0) = u(t, u_0) \,,$$

where  $u(t, u_0) = (w(t), v(t), \theta(t))$  is the solution of the system (11)-(17) with initial conditions  $u_0 = (w_0, w_1, \theta_0)$ , and X is the associated solution space.



Fig. 12

**Proposition 1** Let  $\varphi$  be a continuous semigroup on a metric space X. If the trajectory through  $x \in X$ ,  $\gamma(x) := \bigcup_{t>0} \varphi^t(x)$  is precompact in X, then  $\varphi^t(x) \to \omega(x)$ ,  $t \to \infty$ . Here  $\omega(x)$  is the

 $\omega$ -limit set of x, which is a compact subset of X. Moreover,  $\omega(x)$  is positively invariant with respect to  $\varphi^t$ .

**Definition 9** A *Lyapunov functional* for  $\varphi^t$  is a mapping  $\Phi : X \to \mathbb{R}$ , such that  $\Phi(\varphi^t(x)) \leq \Phi(x)$ , for any  $x \in X, t \geq 0$ .

**Theorem 8** (Dafermos) Let the conditions of Proposition1 hold, and let  $\Phi$  be a Lyapunov functional for  $\varphi$ . Then  $\Phi$  is constant on  $\omega(x)$ .

## 12 Use of the invariance principle for the microwave heating problem

Introduce the Banach space  $X = H_0^1(0, 1) \times L^2(0, 1) \times L^2(0, 1)$ , with the norm

$$\|(w,v,\theta)\|_X^2 = \|w\|_{H^1_0}^2 + \|v\|_{L^2}^2 + \|\theta\|_{L^2}^2.$$

For any  $(w, v, \theta) \in X$  we introduce the functional

$$\Phi(w, v, \theta) = \int_{0}^{1} (w_x^2 + 2\lambda wv + v^2 + a\theta^2) dx, \qquad (18)$$

where  $\lambda$  and a are some positive constants.

(A11) There exist positive parameters  $\lambda$  and a such that for some  $\epsilon > 0$ 

- 1)  $0 < \lambda < 1;$
- 2)  $\frac{1}{2}\sigma_0\epsilon 1 < 0;$
- 3)  $\lambda(\frac{1}{2\epsilon}\sigma_1+1) + ak_1^{-1}\sigma_2 \sigma_0 < 0 \ (\sigma_0, \sigma_1, \sigma_2 \text{ and } k_1 \text{ from the conditions (A9)}).$

**Lemma 1** Under the conditions (A9) – (A11) the functional  $\Phi$  of the system (11)-(17) has the following properties:

1) There exist constants  $C_1, C_2 > 0$ , such that

$$C_1(\|w_x\|_{L^2}^2 + \|v\|_{L^2}^2 + \|\theta\|_{L^2}^2) \le \Phi(w, v, \theta)$$
(19)

$$\leq C_2(\|w_x\|_{L^2}^2 + \|v\|_{L^2}^2 + \|\theta\|_{L^2}^2), \qquad (20)$$

for  $(w, v, \theta) \in X$ .

2) There exists a constant  $C_3 > 0$ , such that

$$\frac{d}{dt}\Phi(t) \le -C_3\Phi(t) \le 0, \quad \text{for all } t \ge t_0, \tag{21}$$

where  $\Phi(t) = \Phi(w(t), v(t), \theta(t))$  and  $(w(t), v(t), \theta(t))$  is some solution of (11)-(17).

Idea of the proof: 1) This inequality follows from the estimate  $||w||_{L^2}^2 \leq ||w_x||_{L^2}^2$  for functions from  $H_0^1(0, 1)$ .

2) Direct computation shows that

$$\begin{aligned} \frac{d}{dt}\Phi(t) &= 2\int_{0}^{1} \left(-\lambda w_{x}^{2} - \lambda \sigma(\theta)wv + (\lambda - \sigma(\theta) + ak'(\theta)^{-1}\sigma(\theta)\theta)v^{2} - ak'(\theta)^{-1}\theta_{x}^{2}\right)dx \leq \\ &- C_{3}\int_{0}^{1} \left(\|w_{x}\|_{L^{2}}^{2} + \|v\|_{L^{2}}^{2} + \|\theta\|_{L^{2}}^{2}\right)dx, \end{aligned}$$

where  $C_3 > 0$  is a constant that exists due to (A11):

1) 
$$\frac{1}{2}\sigma_0\epsilon - 1 < 0;$$

2) 
$$\lambda(\frac{1}{2\epsilon}\sigma_1 + 1) + ak_1^{-1}\sigma_2 - \sigma_0 < 0$$
 or  $\lambda(\frac{1}{2\epsilon}\sigma_1 + 1) + a\sigma_2\sigma_1 - \sigma_0 < 0.$ 

**Theorem 9** If the conditions of the Lemma hold, then any solution of the system (11)-(17) tends in X to the zero solution as  $t \to \infty$ .

**Proof** Define  $\Phi$  as  $\Phi(t) = \Phi(w(t), v(t), \theta(t))$ , where  $(w(t), v(t), \theta(t))$  is some solution of (11)-(17).

Statement 2) of the lemma:

$$\frac{d}{dt}\Phi(t) \le -C_3\Phi(t), \quad t \ge t_0.$$
(22)

Continuity of the functional and Gronwall's inequality:  $\lim_{t\to\infty} \Phi(t) = 0.$ 

#### **13 Numerical experiments**

Initial conditions:  $w_0(x) = 0$ ,  $w_1(x) = p \sin(\pi x)$  and  $\theta_0(x) = p \sin(\pi x)$  for solution components  $w^p(x,t)$  and  $\theta^p(x,t)$ , where p is taken from [-0.5, 0.5].



**Fig. 13** Solution component  $\theta^p(x, t)$ .



**Fig. 14** Solution component  $w^p(x, t)$ .

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14 The two-phase problem

We consider a coefficient of the form



Fig. 15 Noncontinuous two-phase coefficient.

## References

[1] V. Reitmann and H. Kantz, *Frequency domain conditions for the existence of almost periodic solutions in evolutionary variational inequalities.* Stochastics and Dynamics, 4 (3), 483 – 499, 2004.

[2] V. Reitmann, *Convergence in evolutionary variational inequalities with hysteresis nonlinearities.* In: Proc. of Equadiff 11, Bratislava, Slovakia, 2005.

[3] V. Reitmann, *Realization theory methods for the stability investigation of nonlinear infinite-dimensional input-output systems.* In: Proc of Equadiff 12, Brno, Czech, 2009, Mathematica Bohemica, 2010.

[4] G.A. Leonov, and V. Reitmann, Absolute observation stability for evolutionary variational inequalities. World Scientific Publishing Co., Scientific Series on Nonlinear Science, Series B, Vol.14, 2010.