Stability analysis for Maxwell's equation with a thermal effect in one space dimension

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1. Introduction

Let (Q,d) be a metric space called the *base space*. The pair $(\{\tau^t\}_{t\in\mathbb{R}},(Q,d))$ where $\tau^t:Q\to Q$ for each $t\in\mathbb{R}$ is called the *base flow* if

$$\tau^{0} = id_{Q},
\tau^{t} \circ \tau^{s} = \tau^{t+s} \quad \forall t, s \in \mathbb{R}.$$
(1)

Let (M, ρ) be an other metric space (phase space).

Definition 1 The pair $\{\{\varphi^t(q,\cdot)\}_{t\in\mathbb{R}_+,q\in Q},(M,\rho)\}$ where $\varphi^t(q,\cdot):M\to M$ for each $t\in\mathbb{R}_+,q\in Q$ is called a cocycle over the base flow $\{\{\tau^t\}_{t\in\mathbb{R}},(Q,d)\}$ if

$$\varphi^{0}(q,\cdot) = id_{M} \ \forall q \in Q,$$

$$\varphi^{t+s}(q,\cdot) = \varphi^{t}(\tau^{s}(q),\varphi^{s}(q,\cdot)) \ \forall q \in Q, \ \forall t,s \in \mathbb{R}_{+}.$$
(2)

For brevity the cocycle $(\{\varphi^t(q,\cdot)\}_{t\in\mathbb{R}_+,q\in Q},(M,\rho))$ over the base flow $(\{\tau^t\}_{t\in\mathbb{R}},(Q,d))$ will be denoted by (φ,τ) .

Define the space $\widetilde{M} = Q \times M$ with the metric

$$\tilde{\rho}((q_1, u_1), (q_2, u_2)) = \max\{d(q_1, q_2), \rho(u_1, u_2)\},$$

 $(q_i, u_i) \in Q \times M, i = 1, 2$

and the family of mappings $S^t: \widetilde{M} \to \widetilde{M}$, $t \in \mathbb{R}_+$, $S^t(q,u) = (\tau^t(q), \varphi^t(q,u))$.

The dynamical system $(\{S^t\}_{t\in\mathbb{R}_+}, (\widetilde{M}, \widetilde{\rho}))$ is called *skew* product flow.

A non-autonomous set $\widehat{\mathcal{C}} = \{\mathcal{C}(q)\}_{q \in Q}$ is a mapping $Q \to 2^M$. A nonautonomous set is called bounded (closed, compact) if for any $q \in Q$ the set $\mathcal{C}(q)$ is bounded (closed, compact) in M.

A bounded non-autonomous set $\widehat{\mathcal{C}}$ is said to be a *globally B-pullback absorbing* set for (φ, τ) if for any $q \in Q$ and any bounded set $\mathcal{B} \subset M$ there exists a $T = T(q, \mathcal{B})$ such that $\varphi^t(\tau^{-t}(q, \mathcal{B})) \subset \mathcal{C}(q)$ for $t \geq T$.

A non-autonomous set $\widehat{\mathcal{C}}$ is called *globally B-pullback* attracting for (φ,τ) if for any $q\in Q$ and any bounded set $\mathcal{B}\subset M$

$$\lim_{t\to+\infty} \operatorname{dist}(\varphi^t(\tau^{-t}(q),\mathcal{B}),\mathcal{C}(q)) = 0,$$

where dist is the Hausdorff semidistance in (M, ρ) .

A non-autonomous set $\widehat{\mathcal{C}}$ is called *invariant (positively invariant)* for (φ,τ) if for any $q\in Q$ and $t\geq 0$ the equality $\varphi^t(q,\mathcal{C}(q))=\mathcal{C}(\tau^t(q))$ (inclusion $\varphi^t(q,\mathcal{C}(q))\subset \mathcal{C}(\tau^t(q))$) holds.

Definition 2 A non-autonomous set is called a global B-pullback attractor for the cocycle (φ, τ) if it is compact, invariant and is globally B-pullback attracting.

For the existence proof of a B-pullback attractor we will use the following theorem [Kloeden-Schmalfuss, 1987]:

Theorem 1 Let the cocycle (φ, τ) have a compact globally B-pullback absorbing set $\widehat{\mathcal{C}} = \{\mathcal{C}(q)\}_{q \in Q}$.

Then (φ, τ) has a unique B-pullback attractor $\widehat{\mathcal{A}} = \{\mathcal{A}(q)\}_{q \in Q}$, where for each $q \in Q$

$$\mathcal{A}(q) = \cap_{t \in \mathbb{R}_+} \overline{\cup_{s > t, s \in \mathbb{R}_+} \varphi^s(\tau^{-s}(q), \mathcal{C}(\tau^{-s}(q)))}.$$

2. Existence of a B-pullback attractor for the 1-dimensional microwave heating problem

To describe the microwave heating process we will consider for any $T<\infty$ a coupled system of Maxwell's and heat transfer equations:

$$\begin{cases} \varepsilon E_t + \sigma E = \operatorname{rot} H, & (x,t) \in \Omega \times (0,T), \\ \mu H_t + \operatorname{rot} E = 0, & (x,t) \in \Omega \times (0,T), \\ b(\theta)_t - \Delta \theta = \sigma |E|^2, & (x,t) \in \Omega \times (0,T), \end{cases}$$

where E(x,t) is the electric field, H(x,t) is the magnetic field, $\theta(x,t)$ is the temperature, $\sigma(x,\theta)$ is the electrical conductivity, $\varepsilon(x,t)$ is the electric permittivity, $\mu(x,t)$ is the magnetic permeability, and

$$b(s) = \begin{cases} b_1(s), & s < \hat{\theta} \\ \left[b_1(\hat{\theta}), b_2(\hat{\theta})\right], & s = \hat{\theta} \\ b_2(s), & s > \hat{\theta}, \end{cases}$$

is some piecewise smooth function with differentiable, monotone increasing functions $b_1(s)$, $b_2(s)$, such that $b_1(\widehat{\theta}) < b_2(\widehat{\theta})$.

The derivation of the 1-dimensional microwave heating problem is given in [H.-M. Yin et al., 2006]:

$$\begin{cases} \psi_t = \zeta, & x \in (0,1), \ t > 0, \\ \zeta_t = \psi_{xx} - \sigma(\theta)\zeta, & x \in (0,1), \ t > 0, \\ b(\theta)_t = \theta_{xx} + \sigma(\theta)\zeta^2, & x \in (0,1), \ t > 0, \\ \psi(0,t) = 0, \ \psi(1,t) = 0, & t > 0, \\ \theta(0,t) = \theta(1,t) = 0, & t > 0, \\ \psi(x,0) = \psi_0(x), & x \in (0,1), \\ \zeta(x,0) = \zeta_0(x), & x \in (0,1), \\ \theta(x,0) = \theta_0(x), & x \in (0,1). \end{cases}$$

Our goal is to show specific asymptotic behaviour of solution's components, such as asymptotic stability.

Consider the initial-boundary problem

$$w_{tt} - w_{xx} + \sigma(\theta) w_t = 0, \quad 0 < x < 1, \quad t > 0 \theta_t - \theta_{xx} = \sigma(\theta) w_t^2, \qquad 0 < x < 1, \quad t > 0$$
(3)

$$w(0,t) = f_1(t), \quad w(1,t) = f_2(t), \quad t > 0 \theta(0,t) = \theta(1,t) = 0, \quad t > 0$$
(4)

$$w(x,0) = w_0(x), \quad w_t(x,0) = w_1(x), \quad 0 < x < 1$$

 $\theta(x,0) = \theta_0(x), \quad 0 < x < 1$
(5)

where $\theta(x,t)$ is the temperature, w(x,t) is the time integral of the nonzero component of the electric field, $\sigma(\theta)$ is the electric conductivity, $f_1(t), f_2(t)$ are the external perturbations of the electric field.

Assumptions:

- **(A1.1)** σ is locally Lipschitz on $(0, +\infty)$;
- **(A1.2)** There exist constants $0 < \sigma_0 \le \sigma_1$ such that $\sigma_0 \le \sigma(z) \le \sigma_1$ for any z > 0;
- (A1.3) σ is monotone decreasing.
 - (A2) $w_0 \in L^2(0,1), w_1 \in L^2(0,1), \theta_0 \in L^2(0,1), \theta_0 \ge 0$ a.e. on (0,1).
 - (A3) f_1, f_2 are $C^2(\mathbb{R})$ and there exists a constant c such that the functions $|f_1'|, |f_2'|, |f_1''|, |f_2''|$ are bounded on \mathbb{R} by c.

Modification of the existence theorem for weak solutions from [H.-M. Yin et al., 2006] for the 1-dimensional case:

Theorem 2 For any T > 0 there exists a global weak solution $(w(x,t), \theta(x,t))$ of the problem (3)-(5) such that $w \in C([0,T]; L^2(0,1)), \ \theta \in L^2(0,T; H^1(0,1)) \cap C([0,T]; L^2(0,1)).$

Additional assumption:

(A4) The weak solution is unique.

Denote $f(x,t) = f_1(t)(1-x) + f_2(t)x$ and $\psi(x,t) = w(x,t) - f(x,t)$ and introduce the system with homogeneous boundary conditions, i.e.

$$\psi_t = \zeta - f_t,$$

$$\zeta_t = \psi_{xx} - \sigma(\theta)\zeta,$$

$$\theta_t = \theta_{xx} + \sigma(\theta)(\psi_t + f_t)^2, \quad 0 < x < 1, \quad t > 0$$
(6)

with initial and boundary conditions

$$\psi(0,t) = \psi(1,t) = 0, \quad \theta(0,t) = \theta(1,t) = 0, \quad t > 0$$
(7)

$$\psi(x,0) = \psi_0(x) = w_0(x) - f(x,0), \quad 0 < x < 1$$

$$\zeta(x,0) = \zeta_0(x) = w_1(x) - f_t(x,0), \quad 0 < x < 1$$

$$\theta(x,0) = \theta_0(x), \quad 0 < x < 1.$$
(8)

Transformed assumption (A2):

(A2')
$$\psi_0 \in H^1_0(0,1), \zeta_0 \in L^2(0,1), \theta_0 \in L^2(0,1), \theta_0 \geq 0$$
 a.e. on $(0,1)$.

Introduction of the cocycle corresponding to the problem (6)-(8)

Define the metric space

$$M = H_0^1(0,1) \times L^2(0,1) \times (L^2(0,1) \cap \{\theta : \theta \ge 0\})$$
 with the norm

$$\|(\psi,\zeta,\theta)\|_M^2 = \|\psi_x\|_{L^2(0,1)}^2 + \|\zeta\|_{L^2(0,1)}^2 + \|\theta\|_{L^1(0,1)}^2.$$

In our situation: $Q = \mathbb{R}$, $\tau^t(s) = t + s$, $\forall t, s \in \mathbb{R}$

$$\varphi^t(s, u_0) = u(t+s, s, u_0),$$

where $u(t, s, u_0) = (\psi(\cdot, t), \zeta(\cdot, t), \theta(\cdot, t))$ is the solution of (6)-(8) such that $u(s, s, u_0) = u_0 = (\psi_0, \zeta_0, \theta_0)$.

From existence and uniqueness of the solution we conclude (I. Ermakov, Y. Kalinin, V. Reitmann, 2011):

Theorem 3 The system (6)-(8) generates a cocycle $(\{\varphi^t(s,\cdot)\}_{t\in\mathbb{R}_+,s\in\mathbb{R}},(M,\|\cdot\|_M))$ over the base flow $(\{\tau^t\}_{t\in\mathbb{R}},\mathbb{R}).$

Proof of the existence of an absorbing set:

- Lyapunov function for the 1st subsystem (damped wave equation);
- monotonicity methods for the 2nd subsystem (heat equation).

Damped wave equation. Consider the initial-boundary problem for the wave equation separately:

$$\psi_{tt} - \psi_{xx} + \sigma(x, t)\psi_t = f_{tt} - \sigma(x, t)f_t, \quad 0 < x < 1, \quad t > s$$
(9)

$$\psi(0,t) = \psi(1,t) = 0, \quad t > s \tag{10}$$

$$\psi(x,s) = \psi_0, \quad \psi_t(x,s) = \psi_1, \quad 0 < x < 1$$
 (11)

where $s \in \mathbb{R}$. Here $\sigma(x,t)$ is a certain function.

Modified assumptions (A1)-(A3):

(A1.2*) There exist constants $0 < \sigma_0 \le \sigma_1$ such that $\sigma_0 \le \sigma(x,t) \le \sigma_1$ for all $x \in (0,1), t \ge 0$.

- (A2*) $\psi_0 \in H_0^1(0,1), \psi_1 \in L^2(0,1).$
- **(A3*)** The function f(x,t) is C^1 in x, C^2 in t and there exists a constant c>0 such that $|f_t|< c$, $|f_{xt}|< c$, $|f_{tt}|< c$ for any $x\in(0,1),t\in\mathbb{R}$.

Under the assumptions **(A1.2*)** - **(A3*)** the problem (9 - 11) has a unique solution $(\psi(\cdot,t),\psi_t(\cdot,t)) \in M_1 = H_0^1(0,1) \times L^2(0,1)$ [R. Temam, 1993]. For $(\psi,\zeta) \in M_1$ define

$$\|(\psi,\zeta)\|_{M_1}^2 = \|\psi_x\|_{L^2(0,1)}^2 + \|\zeta\|_{L^2(0,1)}^2$$

Write equation (9) as first order system

$$\psi_t = \zeta - f_t,
\zeta_t = \psi_{xx} - \sigma(x, t)\zeta$$
(12)

with boundary and initial conditions

$$\psi(0,t) = \psi(1,t) = 0, \quad t > s \tag{13}$$

$$\psi(x,s) = \psi_0(x), \zeta(x,s) = \zeta_0(x), \quad 0 < x < 1$$
 (14)

Proposition 1 For any t > 0 there exist T > 0, c > 0 such that $\|(\psi(\cdot,t;s),\zeta(\cdot,t;s))\|_{M_1} < c$ for any $s \le t - T$.

Idea of the proof: Lyapunov functional on M_1

$$V(\psi, \zeta) = \|\psi_x\|^2 + 2\lambda(\psi, \zeta) + \|\zeta\|^2$$

where $\lambda > 0$ is a parameter. $\|\cdot\|$ and (\cdot, \cdot) are in $L^2(0,1)$.

Denote $V(t) = V(\psi(\cdot, t), \zeta(\cdot, t))$.

We prove that there exist $\delta > 0$, $c_1 > 0$ and $c_2 \in \mathbb{R}$ such that

$$\frac{d}{dt}V\left(t\right) \leq -\delta V\left(t\right) + c_{1},$$

$$V(t) \le e^{-\delta(t-s)}V(s) + c_2,$$

for any $t, s, t \ge s$.

The nonlinear heat equation (2nd equation of (6))

General setting (A.A.Pankov, 1983):

Suppose that $E \subset H \subset E'$ is a Gelfand triple, i.e. $(E, \|\cdot\|_E)$ is a reflexive Banach space, H is a Hilbert space, E' is the space dual to E, E is continuously and densely embedded into H.

Suppose that $A(t): E \to E'$ is a family of operators and $\tilde{f}: \mathbb{R} \to E'$ is a measurable function . The operator $A(t): E \to E'$ is monotone, i.e.

$$(A(t)u - A(t)v, u - v) \ge 0, \ \forall u, v \in E.$$

Here (\cdot, \cdot) is the duality pairing on $E \times E'$, coinsiding on $E \times E$ with the scalar product in H.

Consider the evolution equation

$$\frac{du}{dt} + A(t)u = \tilde{f}(t). \tag{15}$$

Suppose that there is an $\alpha > 0$ such that

$$(A(t)u - A(t)v, u - v) \ge \alpha \|u - v\|^2 \ \forall u, v \in E.$$
 (16)

Define the following function spaces:

- $C_b(\mathbb{R}, E)$ is the set of continuous functions $f: \mathbb{R} \to E$, for which $\sup_{t \in \mathbb{R}} \|f(t)\|_E$ is finite.
- $BS^p(\mathbb{R}, E), 1 \leq p < \infty$ is the subspace in $L^p_{loc}(\mathbb{R}, E)$, consisting of functions with finite norm

$$\|f\|_{S^p}^p = \sup_{t \in \mathbb{R}} \left(\int_t^{t+1} \|f(s)\|_E^p ds \right).$$

Consider the heat equation in the form

$$\theta_t - \theta_{xx} = \sigma(\theta)g(x,t).$$

Suppose that $g(x,t) \ge 0$ is measurable and uniformly bounded in t. We have $g(x,t) = (\psi_t(x,t) + f_t(x,t))^2$. For σ the assumptions (A1.1)-(A1.3) hold.

 $\sigma(\theta) = \sigma_0 + \tilde{\sigma}(\theta)$, where σ_0 is from (A1.2). We get

$$\theta_t - \theta_{xx} - \tilde{\sigma}(\theta)g(x,t) = \sigma_0 g(x,t), \quad 0 < x < 1, \quad t > s$$
(17)

$$\theta(0,t) = \theta(1,t) = 0, \quad t > s$$
 (18)

$$\theta(x,s) = \theta_0(x), \quad 0 < x < 1.$$
 (19)

The initial-boundary problem (17)-(19) generates an evolution equation (15), where

$$A(t)u = -u'' - g(x,t)\tilde{\sigma}(u)$$
 for $u \in E$ and $\tilde{f}(t) = \sigma_0 g(\cdot,t)$.

In our situation we have $E=H_0^1(0,1)$ and $H=L^2(0,1)$. Check condition (16). Let $u,v\in H_0^1(0,1)$, $\eta=u-v$. Then

$$(A(t)u - A(t)v, u - v) = (-\eta'', \eta) + (g(\cdot, t)(\tilde{\sigma}(v) - \tilde{\sigma}(u)), \eta) =$$
$$= (\eta', \eta') + (g(\cdot, t)(\tilde{\sigma}(v) - \tilde{\sigma}(u)), \eta) \ge \|\eta\|^2.$$

A.A. Pankov, 1983:

- 1. The Cauchy problem for equation (15) has a unique solution $u \in BS^2(\mathbb{R}, H_0^1(0,1)) \cap C_b(\mathbb{R}, L^2(0,1))$. For the equation (17) this means that there exists a constant c_1 such that $\|\theta(\cdot,t;s)\| \leq c_1$ for any $t,s \in \mathbb{R}$, $s \leq t$.
- 2. We have the estimate

$$\|\theta_1(\cdot,t;s) - \theta_2(\cdot,t;s)\| \le e^{-c_2(t-s)} \|\theta_{01} - \theta_{02}\|,$$
 (20)

for t > s, where $\theta_i(x, t; s)$ is the solution of (17) with initial data θ_{0i} and initial time s.

The constant c_1 does not depend on initial data:

$$\theta(x,t;s) = \int_0^1 G(x,y;t,s)\theta_0(y)dy + \int_s^t \int_0^1 G(x,y;t,r)g(y,r)drdy,$$

where G(x, y; t, r) is the corresponding Green's function which satisfies

$$|G(x,y;t,s)| \le \frac{c_3}{\sqrt{t-s}}.$$

The influence of initial data tends to zero for $t \to \infty$.

Make the initial time s tend to $-\infty$, which corresponds to the time shift in g(x,t).

Proposition 2 Let $\theta(\cdot,t;s)$ be the solution of (17)-(19). There exists a constant c such that for all t and $s \le t$ the inequality $\|\theta(\cdot,t;s)\| \le c$ holds where c does not depend on θ_0 .

From uniform boundedness in s of solutions of the wave equation and the heat equation we obtain

Theorem 4 The cocycle (φ, τ) generated by problem (6) - (8) has a globally B-pullback absorbing set.

Applying the Kloeden-Schmalfuss Theorem, we get

Theorem 5 The cocycle (φ, τ) generated by problem (6) - (8) has a global B-pullback attractor.

3. Determining functionals for cocycles

Physical meaning: Asymptotically finite-dimensional dynamics

C. Foias, G. Prodi, 1967

O. Ladyzhenskaya, 1975

I.D. Chueshov, 1998

I.D. Chueshov, J. Duan, B. Schmalfuss, 2001.

If the system has a global attractor, such functionals can give an approximation of the attractor.

Let $(\{S^t\}_{t\in\mathbb{R}_+}, (E, \|\cdot\|))$ be a dynamical system on Banach space $(E, \|\cdot\|)$.

Definition 3 The set $\{l_j\}_{j=1}^N$ of linear continuous functionals on E is called a set of asymptotically determining functionals for the dynamical system $(\{S^t\}_{t\in\mathbb{R}_+},(E,\|\cdot\|))$ if for any $u_1,u_2\in E$ the condition

$$\lim_{t \to +\infty} |l_j(S^t(u_1)) - l_j(S^t(u_2))| = 0, \ j = 1, ..., N$$

implies

$$\lim_{t \to +\infty} ||S^t(u_1) - S^t(u_2)|| = 0.$$

Introduce the determining modes which are important examples of determining functionals.

Definition 4 The determining modes for the dynamical system $(\{S^t\}_{t\in\mathbb{R}_+}, (H, (\cdot, \cdot)))$ on a Hilbert phase space $(H, (\cdot, \cdot))$ are determining functionals $l_j(\cdot) = (\cdot, e_j)$ where $\{e_j\}_1^N$ are some elements of H.

The notion of pullback-asymptotically determining functionals for processes was introduced in [J.A. Langa, 2003]. We give a generalization for cocycles.

Definition 5 The set $\{l_j\}_{j=1}^N$ of linear continuous functionals on Banach space $(M,\|\cdot\|)$ is called a set of pullback-asymptotically determining functionals for the cocycle $(\{\varphi^t(q,\cdot)\}_{q\in Q,t\in\mathbb{R}_+},(M,\|\cdot\|))$ over the base flow $(\{\tau^t\}_{t\in\mathbb{R}},(Q,d))$ if the condition

$$\lim_{t\to+\infty} \left| l_j(\varphi^t(\tau^{-t}(q),u_1)) - l_j(\varphi^t(\tau^{-t}(q),u_2)) \right| = 0$$

for any $q \in Q$, $u_1, u_2 \in M$, j = 1, ..., N implies

$$\lim_{t\to+\infty} \left\| \varphi^t(\tau^{-t}(q),u_1) - \varphi^t(\tau^{-t}(q),u_2) \right\| = 0.$$

Let (φ, τ) be a cocycle on a Hilbert phase space H, π_1 be the projector from H onto a finite-dimensional subspace of H and π_2 be its complement.

Assumptions:

- **(H1)** The non-autonomous set $\{C(q)\}_{q\in Q}$ is positively invariant for (φ, τ) .
- **(H2)** For any $q \in Q$ there exists $\delta = \delta(q) \in (0,1)$ such that for all $s \geq 1, u, v \in \mathcal{C}(\tau^{-s}(q))$

$$\|\pi_2(\varphi^1(\tau^{-s}(q),u)-\varphi^1(\tau^{-s}(q),v))\| \leq \delta(q)\|u-v\|.$$

Let $a_1, a_2 : Q \to H$ be mappings such that $a_i(q) \in C(q)$ for any $q \in Q$.

(H3) For any $\varepsilon > 0$, $t \ge 0$ there exists an $L = L(\varepsilon) \in \mathbb{N}$ such that for any $q \in Q$

$$\delta(q)^{2L} \left\| \varphi^{t-L}(q, a_1(q)) - \varphi^{t-L}(q, a_2(q)) \right\|^2 < \varepsilon,$$
 and $L(\varepsilon) \to \infty$ if $\varepsilon \to 0$.

The next theorem [I. Ermakov, Y. Kalinin, V. Reitmann, 2011] is a generalization of Theorem 14 from [J.A. Langa, 2003]

Theorem 6 Let the assumptions (H1)-(H3) hold and suppose that there exists a $\beta > 0$ such that for any $q \in Q$

$$\lim_{t\to+\infty} \left\| \pi_1(\varphi^t(\tau^{-t}(q),a_1(q)) - \varphi^t(\tau^{-t}(q),a_2(q))) \right\| \leq \beta.$$

Then

$$\lim_{t \to +\infty} \| \varphi^t(\tau^{-t}(q), a_1(q)) - \varphi^t(\tau^{-t}(q), a_2(q)) \| \le \beta. \quad (21)$$

Corollary 1 Let there exist a $\beta > 0$ such that for all $q \in Q, u, v \in H$

$$\lim_{t\to+\infty} \left\| \pi_1(\varphi^t(\tau^{-t}(q),u) - \varphi^t(\tau^{-t}(q),v)) \right\| \leq \beta.$$

Then

$$\lim_{t \to +\infty} \left\| \varphi^t(\tau^{-t}(q), u) - \varphi^t(\tau^{-t}(q), v) \right\| \le \beta.$$

This corollary gives the existence of pullbackasymptotically determining modes for a cocycle.

Now consider cocycles of a special type. Such cocycles are generated by the microwave heating problem.

Let (φ, τ) be a cocycle with phase space $E = E_1 \times E_2$ where E_1 is a Hilbert and E_2 is a Banach space, respectively.

 φ has the form (φ_1, φ_2) , i.e.

$$\varphi_1: \mathbb{R}_+ \times Q \times E_1 \times E_2 \to E_1$$
,

$$\varphi_2: \mathbb{R}_+ \times Q \times E_1 \times E_2 \to E_2.$$

Let π_1 be the projector from E_1 onto a finite-dimensional subspace of E_1 , π_2 be its orthogonal complement.

Let $a_1, a_2 : Q \to E$ be mappings such that $a_i(q) \in \mathcal{C}(q)$ for any $q \in Q$.

Modify assumptions (H2) and (H3) so that they hold for φ_1 instead of φ :

- (H2*) For any $q \in Q$ there exists a $\delta = \delta(q) \in (0,1)$ such that for all $s \geq 1, u, v \in \mathcal{C}(\tau^{-s}(q))$ we have $\left\|\pi_2(\varphi_1^1(\tau^{-s}(q), u) \varphi_1^1(\tau^{-s}(q), v))\right\|_{E_1} \leq \delta(q) \|u v\|_E.$
- **(H3*)** For any $\varepsilon>0$, $t\geq 0$ there exists an $L=L(\varepsilon)\in\mathbb{N}$ such that for any $q\in Q$

$$\delta(q)^{2L} \left\| \varphi_1^{t-L}(q, a_1(q)) - \varphi_1^{t-L}(q, a_2(q)) \right\|_{E_1}^2 < \varepsilon,$$
 and $L(\varepsilon) \to \infty$ if $\varepsilon \to 0$.

Theorem 7 Suppose that

- 1) The assumptions (H1),(H2*),(H3*) hold for the cocycle (φ, τ) .
- 2) The estimate

$$\begin{split} \left\| \varphi_2^t(\tau^{-t}(q), u_1, u_2) - \varphi_2^t(\tau^{-t}(q), v_1, v_2) \right\|_{E_2} &\leq e^{-ct} \left\| u_2 - v_2 \right\|_{E_2} \\ \text{holds with some constant } c > 0 \text{ for any } t > 0, \ u_1, v_1 \in E_1, \ u_2, v_2 \in E_2. \end{split}$$

3) There exists $a \beta > 0$ such that for any $q \in Q$ $\lim_{t \to +\infty} \left\| \pi_1(\varphi_1^t(\tau^{-t}(q), a_1(q)) - \varphi_1^t(\tau^{-t}(q), a_2(q))) \right\|_{E_1} \leq \beta.$

Then

$$\lim_{t \to +\infty} \| \varphi^t(\tau^{-t}(q), a_1(q)) - \varphi^t(\tau^{-t}(q), a_2(q)) \|_E \le \beta.$$
(22)

4. Numerical experiments

Consider the family of initial conditions of the form $\psi_0(0) = 0$, $\psi_0(x) = p\sin(\pi x)$ and $\theta_0(x) = p\sin(\pi x)$, where p is taken from [0,1]. $\sigma(z) = 2 + \sin(z)$.

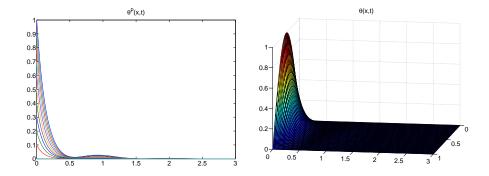


Figure 1 Solution component $\theta^p(x,t)$.

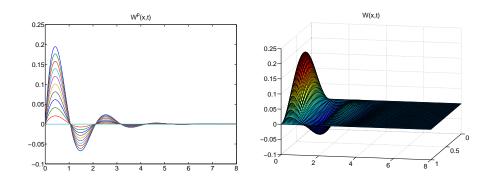


Figure 2 Solution component $\psi^p(x,t)$.

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